ASC Report No. 25/2008

On the Convergence of Adaptive Non-Conforming Finite Element Methods

Christoph Ortner, Dirk Praetorius

Institute for Analysis and Scientific Computing
Vienna University of Technology — TU Wien
www.asc.tuwien.ac.at ISBN 978-3-902627-00-1

Most recent ASC Reports

24/2008	*Christoph Ortner, Dirk Praetorius* A Non-Conforming Finite Element Method for Convex Variational Problems
23/2008	*Samuel Ferraz-Leite, Dirk Praetorius* Adaptive Boundary Element Methods Based on Accurate a Posteriori Error Estimation
22/2008	*Jeff Cash, Georg Kitzhofer, Othmar Koch, Gerald Moore, Ewa Weinmüller* Numerical Solutions of singular Two Point BVPs
21/2008	*Samuel Ferraz-Leite, Dirk Praetorius* A Posteriori Fehlerschätzer für die Symmsche Integralgleichung in 3D
20/2008	*Bertram Düring* Asset Pricing under Information with Stochastic Volatility
19/2008	*Stefan Funken, Dirk Praetorius, Philipp Wissgott* Efficient Implementation of Adaptive P1-FEM in MATLAB
18/2008	*Bertram Düring, Guiseppe Toscani* International and Domestic Trading and Wealth Distribution
17/2008	*Vyacheslav Pivovarchik, Harald Woracek* Sums of Nevanlinna Functions and Differential Equations on Star-Shaped Graphs
16/2008	*Bertram Düring, Daniel Matthes, Guiseppe Toscani* Kinetic Equations Modelling Wealth Redistribution: A Comparison of Approaches
15/2008	*Jens Markus Melenk, Stefan Sauter* Convergence Analysis for Finite Element Discretizations of the Helmholtz Equation. Part I: the Full Space Problem

Institute for Analysis and Scientific Computing
Vienna University of Technology
Wiedner Hauptstraße 8–10
1040 Wien, Austria

E-Mail: admin@asc.tuwien.ac.at
WWW: http://www.asc.tuwien.ac.at
FAX: +43-1-58801-10196

ISBN 978-3-902627-00-1

ON THE CONVERGENCE OF ADAPTIVE NON-CONFORMING FINITE ELEMENT METHODS

C. ORTNER & D. PRAETORIUS

Abstract. We formulate and analyze an adaptive non-conforming finite-element method for the solution of convex variational problems. The class of minimization problems we admit includes highly singular problems for which no Euler–Lagrange equation (or inequality) is available. As a consequence, our arguments only use the structure of the energy functional. We are nevertheless able to prove convergence of an adaptive algorithm, despite using refinement indicators which are not reliable error indicators.

1. Introduction

We present a new convergence proof for adaptive non-conforming finite element methods which is applicable to a wide class of convex variational problems.

For fixed $n \geq 2$ and $m \geq 1$, let Ω be a bounded polygonal domain in $\mathbb{R}^n$ and let $W : \mathbb{R}^{m\times n} \to [0, +\infty]$ be a *convex* stored energy density which satisfies a p-growth condition from below only, i.e.,

$$W(\xi) \geq C(|\xi|^p - 1) \quad \text{for all } \xi \in \mathbb{R}^{m\times n} \text{ and some } p > 1. \tag{1}$$

For a fixed dead load $f \in \mathrm{L}^{p'}(\Omega)$, where $1/p' + 1/p = 1$, we define the energy functional $\mathcal{J} : \mathrm{W}^{1,p}(\Omega)^m \to \mathbb{R} \cup \{+\infty\}$ by

$$\mathcal{J}(v) = \int_\Omega \big(W(\nabla v) - f \cdot v\big)\,\mathrm{d}x. \tag{2}$$

Given some $g \in \mathrm{W}^{1,p}(\Omega)^m$ with $\mathcal{J}(g) < +\infty$ and sets $\Gamma^{(i)} \subset \partial\Omega$, $i = 1, \dots, m$, with positive surface measure $|\Gamma^{(i)}| > 0$, we define the *admissible set*

$$\mathcal{A} = \big\{v \in \mathrm{W}^{1,p}(\Omega)^m \,:\, v^{(i)} = g^{(i)} \text{ on } \Gamma^{(i)}, i = 1, \dots, m\big\} \neq \emptyset; \tag{3}$$

here and throughout we use superscripts to denote components of a vector-valued function. We note that admissible functions satisfy the Poincaré-type inequality

$$\|v - g\|_{\mathrm{L}^p(\Omega)} \leq C_p \|\nabla v - \nabla g\|_{\mathrm{L}^p(\Omega)} \quad \text{for all } v \in \mathcal{A}, \tag{4}$$

where C_p is some fixed constant.

Date: September 29, 2008.

2000 *Mathematics Subject Classification.* 49J45, 65N30, 49M25.

Key words and phrases. Calculus of Variations, Crouzeix–Raviart FEM, Convergence of AFEM, Lavrentiev Phenomenon.

The authors acknowledge the kind support and generous hospitality of the Hausdorff Centre for Mathematics (University of Bonn) where this work was initiated.

In this paper, we analyze the numerical solution of the minimization problem

$$u \in \operatorname{argmin} \mathcal{J}(\mathcal{A}). \tag{5}$$

by means of an adaptive non-conforming finite element method.

Under the conditions we imposed, the functional $\mathcal{J}$ may be non-differentiable at a solution, even if W itself is smooth [3], and hence the associated Euler–Lagrange equation is unavailable to us. In this situtation, the method of choice for the analysis of (5) (and its discretization) is the *direct method of the calculus of variations* [15], whose application immediately gives the following basic existence result. Since it is helpful to understand our subsequent analysis, we give a brief outline of the proof. We refer to Dacorogna's monograph [15] (in particular Theorem 4.1) for further details.

Proposition 1. *There exists at least one solution of* (5).

Proof. For an admissible function $v \in \mathcal{A}$, the Poincaré inequality (4) implies

$$\|v\|_{L^p(\Omega)} \lesssim \|\nabla(v-g)\|_{L^p(\Omega)} + \|g\|_{L^p(\Omega)} \lesssim \|\nabla v\|_{L^p(\Omega)} + \|g\|_{W^{1,p}(\Omega)},$$

where $a \lesssim b$ abbreviates $a \le C\, b$ with a generic constant $C > 0$. From the p-growth condition (1) and Young's inequality, we infer

$$\mathcal{J}(v) \gtrsim \|\nabla v\|_{L^p(\Omega)}^p - \|f\|_{L^{p'}(\Omega)} \|v\|_{L^p(\Omega)} - |\Omega| \gtrsim \|\nabla v\|_{L^p(\Omega)}^p - C(f, g, \Omega)$$

with an additive constant $C(f, g, \Omega) \ge 0$. Applying the Poincaré-type inequality (4) again, we find that $\mathcal{J}$ is coercive, i.e.,

$$\|v\|_{W^{1,p}(\Omega)}^p \lesssim \mathcal{J}(v) + 1 \quad \text{for all } v \in \mathcal{A}. \tag{6}$$

Suppose now that $(u_\ell)_{\ell\in\mathbb{N}} \subset \mathcal{A}$ is a minimizing sequence for $\mathcal{J}$, i.e., $\mathcal{J}(u_\ell) \to \inf \mathcal{J}(\mathcal{A})$ as $\ell \to \infty$. By (6) and reflexivity of $W^{1,p}(\Omega)^m$, we may assume that $(u_\ell)_{\ell\in\mathbb{N}}$ converges weakly to some limit u in $W^{1,p}(\Omega)^m$. Since $\mathcal{A}$ is convex and closed, it is weakly closed, and hence it follows that $u \in \mathcal{A}$. Finally, convexity and non-negativity of W imply the weak lower semicontinuity of $\mathcal{J}$, i.e.,

$$\mathcal{J}(u) \le \liminf_{\ell\to\infty} \mathcal{J}(u_\ell),$$

cf. [15, Theorem 3.4]. Thus, we conclude that $\mathcal{J}(u) = \inf \mathcal{J}(\mathcal{A})$, i.e., that $u \in \operatorname{argmin} \mathcal{J}(\mathcal{A})$. □

The class of problems which is included in our analysis is surprisingly general. It includes not only standard variational model problems such as the Dirichlet problem [17, 19, 10, 13] or the p-Laplacian [22, 16] or even convex problems with control of stresses [4, 9] (all of these works require uniform p-growth from below *and* above), but is in fact suited for any convex functional of the type (2) with convex constraints. Even our p-growth condition from below can be relaxed to some extent [21]. For simplicity we have restricted ourselves to Dirichlet conditions, but this constraint is easily lifted as well.

One of the most important features of our analysis is that we do not require $\mathcal{J}$ to be continuous in the strong topology of $W^{1,p}(\Omega)^m$ and hence, using ideas from [21], we are even able to overcome the *Lavrentiev gap phenomenon.* If $g \in W^{1,\infty}(\Omega)^m$, then

$\mathcal{A}_\infty := \mathcal{A} \cap \mathrm{W}^{1,\infty}(\Omega)^m$ is non-empty, and we say that (5) exhibits the Lavrentiev gap phenomenon if

$$\inf \mathcal{J}(\mathcal{A}_\infty) > \min \mathcal{J}(\mathcal{A}). \tag{7}$$

Foss, Hrusa and Mizel [18] have shown that this effect can indeed occur under the conditions we have posed. Note that, since (most) conforming finite element functions are Lipschitz continuous, non-conformity of the numerical method is essential for treating problems in that class. Furthermore, since the Lavrentiev effect is closely linked to singularities in the solution of nonlinear variational problem, adaptive solution techniques are particularly important. We refer to [3, 12] for overviews of this fascinating subject.

To the best of our knowledge, three classes of convergence proofs for adaptive finite element methods for linear problems exist to date. The first idea [17, 19] used the so-called inner node property from which a lower bound on the discrete error in terms of the estimator can be derived which implies, up to oscillation terms, the reduction of the error at each step of the adaptive algorithm. A further step was taken in [20] by circumventing the inner node property at a given step but still requiring that it should be obtained after a fixed number of refinements. One particular aspect of their strategy of proof is also used in our approach. The first proof which completely circumvented the use of lower bounds but only requires reliability of the error estimator, was recently found by Cascon, Kreuzer, Nochetto, and Siebert [13]. All these proofs rely heavily on the use of Galerkin orthogonality, which is unavailable to us. Extensions of this convergence analysis to the non-linear Laplacian, or more general convex variational problem can be found in [4, 9, 16, 22]. These works require the use of some additional structure of the problems but are still heavily based on the analysis for the Dirichlet problem.

Proofs of the convergence of adaptive non-conforming and mixed finite element methods can be found, for example, in [11, 10]. The analysis contained therein is mostly adapted from the conforming case, the crucial modification being a control of non-conformity which leads to so-called quasi-Galerkin orthogonality relations.

In particular, all previous proofs require the use of Euler–Lagrange equations which are not available for our problem class. We therefore have to use a completely different approach. Our convergence argument is strictly tailored to non-conforming finite element methods and does not apply in an obvious way to the conforming case. In addition to the lack of Euler–Lagrange equations and Galerkin orthogonality in our problem, it is also interesting to note that we obtain convergence of our adaptive algorithm even though the refinement indicators are not reliable error bounds, but rather *conformity indicators*.

The paper is organized as follows. In Section 2, we fix the notation and state some auxiliary results. In Section 3, we provide sufficient conditions for the convergence of the Crouzeix–Raviart finite element method. This analysis motivates the definition of convergence indicators which we discuss in some detail in Section 3.2. Finally, in Section 3.3 we formulate an adaptive algorithm and prove its convergence. To conclude, we present several numerical experiments in Section 4.

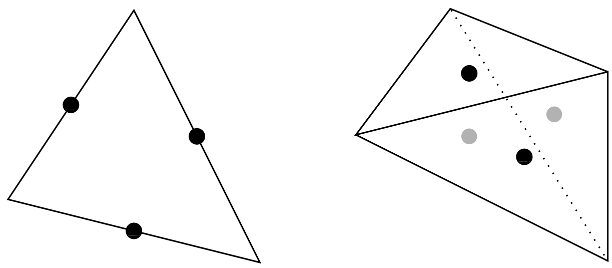

FIGURE 1. Crouzeix–Raviart elements in two and three dimensions.

2. Preliminaries

2.1. Function Spaces. Let A be an open subset of $\mathbb{R}^n$. We use $\mathrm{L}^p(A)$ and $\mathrm{W}^{1,p}(A)$ to denote the standard Lebesgue and Sobolev spaces and equip them with their usual norms. The space of distributions is denoted by $\mathcal{D}'(A)$ [2]. The distributional gradient operator is denoted D, while the weak gradient operator is denoted ∇. The space of k-times continuously differentiable functions with compact support in A is denoted $\mathrm{C}_0^k(A)$.

2.2. Triangulation of Ω. For every $\ell \in \mathbb{N}$, we assume that $\mathcal{T}_\ell$ is a regular triangulation of $\bar\Omega$ into closed simplices $T \in \mathcal{T}_\ell$. In particular, $\mathcal{T}_\ell$ has no hanging nodes in 2D, no hanging nodes or edges in 3D, and so forth. Let $\mathcal{E}_\ell$ denote the collection of $n-1$ dimensional faces of elements and $\mathcal{E}_\ell^{\rm int}$ the collection of interior faces. Let $\mathcal{N}_\ell$ denote the set of vertices and $\mathcal{N}_\ell^{\rm nc}$ denote the set of barycenters of the faces. We assume throughout that, up to surface measure zero, the sets $\Gamma^{(i)}$ are the union of faces $\mathcal{E}_\ell^{d,i} \subset \mathcal{E}_\ell$, $i = 1, \dots, m$, and we set $\mathcal{E}_\ell^{{\rm int},i} = \mathcal{E}_\ell \setminus \mathcal{E}_\ell^{d,i}$.

For each element $T \in \mathcal{T}_\ell$, we set $h_T = \mathrm{diam}(T)$. For each face $E \in \mathcal{E}$, we set $h_E = \mathrm{diam}(E)$. The mesh-size function $h_\ell : \Omega \to \mathbb{R}_{>0}$ is then almost everywhere defined by $h_\ell(x) = h_T$ for x in the interior of an element $T \in \mathcal{T}_\ell$ and $h_\ell(x) = h_E$ for x in the relative interior of a face $E \in \mathcal{E}_\ell$.

The shape regularity constant $\sigma(\mathcal{T}_\ell)$ is the smallest number $C > 0$ such that

$$C^{-1}\, h_T^n \le |T| \le h_T^n, \quad C^{-1}\, h_E^{n-1} \le |E| \le h_E^{n-1} \quad \text{and} \quad h_E \le h_T \le C\, h_E$$

for all elements $T \in \mathcal{T}_\ell$ and faces $E \in \mathcal{E}_\ell$ with $E \subset \partial T$. A family $(\mathcal{T}_\ell)_{\ell\in\mathbb{N}}$ of regular triangulations is uniformly shape-regular, if $\sup\limits_{\ell\in\mathbb{N}} \sigma(\mathcal{T}_\ell) < \infty$.

2.3. Finite Element Spaces. We briefly describe the Crouzeix–Raviart finite element space used to discretize (5); see [5, 7, 14] for further details.

The space of all $\mathcal{T}_\ell$-piecewise affine functions is denoted by

$$\mathrm{P}_1(\mathcal{T}_\ell) = \big\{ v \in \mathrm{L}^1(\Omega) \,:\, v|_T \text{ is affine for all } T \in \mathcal{T}_\ell \big\}.$$

$\mathrm{CR}(\mathcal{T}_\ell)$ denotes the Crouzeix–Raviart finite element space,

$$\mathrm{CR}(\mathcal{T}_\ell) = \big\{ v_\ell \in \mathrm{P}_1(\mathcal{T}) \,:\, v_\ell \text{ is continuous in all face barycentres } z \in \mathcal{N}_\ell^{\rm nc} \big\}.$$

Diagrams of the Crouzeix–Raviart elements in two and three dimensions are shown in Figure 1.

For each barycenter $z \in \mathcal{N}_\ell^{\mathrm{nc}}$, let $E_z \in \mathcal{E}_\ell$ be the unique face which contains z. The interpolation operator $\Pi_\ell : \mathrm{W}^{1,1}(\Omega)^m \to \mathrm{CR}(\mathcal{T}_\ell)^m$ is defined via

$$\Pi_\ell v(z) = |E_z|^{-1} \int_{E_z} v \,\mathrm{d}s \quad \text{for all } z \in \mathcal{N}_\ell^{\mathrm{nc}}$$

and thus satisfies $\int_E \Pi_\ell v \,\mathrm{d}s = \int_E v \,\mathrm{d}s$ for all faces $E \in \mathcal{E}_\ell$. This operator was first used by Crouzeix and Raviart [14]. We summarize its most important properties for our purpose in the following lemma. It is worth noting that neither the stability of Π_ℓ, nor the interpolation error estimate depends on the shape regularity $\sigma(\mathcal{T}_\ell)$.

Lemma 2 (Properties of Crouzeix–Raviart Interpolant). *Let $v \in \mathrm{W}^{1,p}(\Omega)^m$ and $T \in \mathcal{T}_\ell$, then Π_ℓ has the first-order approximation property*

$$\|v - \Pi_\ell v\|_{\mathrm{L}^p(T)} \le C_{\mathrm{apx}} h_T \,\|\nabla v\|_{\mathrm{L}^p(T)} \tag{8}$$

with $C_{\mathrm{apx}} = 1 + 2/n \le 2$. Furthermore, it satisfies the mean value property

$$|T|^{-1} \int_T \nabla v \,\mathrm{d}x = |T|^{-1} \int_T \nabla(\Pi_\ell v) \,\mathrm{d}x = \nabla(\Pi_\ell v)|_T, \tag{9}$$

which implies the stability estimate

$$\|\nabla(\Pi_\ell v)\|_{\mathrm{L}^p(T)} \le \|\nabla v\|_{\mathrm{L}^p(T)}. \tag{10}$$

Proof. We use integration by parts on the element T. Keeping in mind that the outer unit normal ν to T is constant on each face, we observe

$$\int_T \frac{\partial}{\partial x_j} v \,\mathrm{d}x = \int_{\partial T} v \nu_j \,\mathrm{d}s = \int_{\partial T} (\Pi_\ell v)\nu_j \,\mathrm{d}s = \int_T \frac{\partial}{\partial x_j} (\Pi_\ell v) \,\mathrm{d}x,$$

which proves (9). In particular, we have

$$\|\nabla(\Pi_\ell v)\|_{\mathrm{L}^p(T)} = |T|^{1/p-1} \Big| \int_T \nabla v \,\mathrm{d}x \Big| \le \|\nabla v\|_{\mathrm{L}^p(T)},$$

which shows (10).

Next, we recall the well-known trace identity

$$\frac{1}{|E|} \int_E w \,\mathrm{d}s = \frac{1}{|T|} \int_T w \,\mathrm{d}x + \frac{1}{n|T|} \int_T (x - z) \cdot \nabla w \,\mathrm{d}x \quad \text{for all } w \in \mathrm{W}^{1,1}(T),$$

where $z \in T \cap \mathcal{N}$ is the vertex opposite to E, i.e. $T = \mathrm{conv}(E \cup \{z\})$. By definition, $w := v - \Pi_\ell v$ satisfies $\int_E w \,\mathrm{d}s = 0$. Therefore, the integral mean $w_T := |T|^{-1} \int_T w \,\mathrm{d}x$ can be estimated by

$$|w_T| = \frac{1}{|T|} \Big| \int_T w \,\mathrm{d}x \Big| \le \frac{h_T}{n|T|} \,\|\nabla w\|_{L^1(T)} \le \frac{h_T |T|^{1/p'-1}}{n} \,\|\nabla w\|_{\mathrm{L}^p(T)},$$

whence $\|w_T\|_{L^p(T)} \le (h_T/n)\,\|\nabla w\|_{L^p(T)}$. Next, we use the Poincaré inequality $\|w - w_T\|_{L^p(T)} \le (h_T/2)\,\|\nabla w\|_{L^p(T)}$ on the convex domain T, cf. [1]. This gives

$$\|w\|_{L^p(T)} \le \|w - w_T\|_{L^p(T)} + \|w_T\|_{L^p(T)} \le \Big(\frac{1}{2} + \frac{1}{n}\Big)\, h_T\,\|\nabla w\|_{L^p(T)}.$$

Hence, we can deduce (8) from (10) and a triangle inequality. □

Since elements of $\mathrm{CR}(\mathcal{T}_\ell)^m$ may be discontinuous over faces $E \in \mathcal{E}_\ell$, we use ∇v_ℓ to denote the $\mathcal{T}_\ell$-elementwise gradient of $v_\ell \in \mathrm{CR}(\mathcal{T}_\ell)^m$. We also require a notation for the jumps $[v_\ell]$ across interior faces. For $E = T^+ \cap T^- \in \mathcal{E}_\ell^{\mathrm{int}}$, we fix the labelling of the elements $T^\pm$ and we let $v_\ell^\pm$ denote the traces from $T^\pm$, and $\nu = \nu_E$ the outer unit normals to T^+. We define the jump across E by

$$[v_\ell] = v_\ell^+ - v_\ell^-.$$

For a boundary face $E \subset \partial\Omega$, $E = \partial\Omega \cap T$, we let $\nu = \nu_E$ be the outer unit normal to Ω, and we define

$$[v_\ell]_i = \begin{cases} v_\ell^{(i)} - g^{(i)}, & \text{if } E \subset \Gamma^{(i)} \\ 0, & \text{otherwise.} \end{cases}$$

With this notation, the distributional gradient reads

$$\langle D v_\ell, \varphi\rangle = -\int_\Omega v_\ell \cdot \mathrm{div}\varphi \,\mathrm{d}x = \int_\Omega \nabla v_\ell : \varphi \,\mathrm{d}x - \int_{\cup \mathcal{E}_\ell^{\mathrm{int}}} ([v_\ell] \otimes \nu) : \varphi \,\mathrm{d}s, \tag{11}$$

for $v_\ell \in \mathrm{P}_1(\mathcal{T}_\ell)^m$ and $\varphi \in \mathrm{C}_0^1(\Omega)^{m\times n}$. The representation formula (11) can be verified by integrating by parts on each element. The symbol $\otimes$ denotes the tensor product $a \otimes b \in \mathbb{R}^{m\times n}$, i.e., $(a \otimes b)_{ij} = a_i b_j$.

3. Adaptive Solution

Let $(\mathcal{T}_\ell)_{\ell\in\mathbb{N}}$ be a uniformly shape-regular family of triangulations of Ω, which will subsequently be generated by an adaptive algorithm. We extend the definition of the energy functional $\mathcal{J}$ to the Crouzeix–Raviart finite element space $\mathrm{CR}(\mathcal{T}_\ell)^m$ by setting

$$\mathcal{J}(v_\ell) = \int_\Omega \Big[W(\nabla v_\ell) - f \cdot v_\ell\Big]\,\mathrm{d}x.$$

We stress, however, that ∇v_ℓ now denotes the $\mathcal{T}_\ell$-piecewise gradient.

Since $\mathrm{CR}(\mathcal{T}_\ell)$ is not a subspace of $\mathrm{W}^{1,p}(\Omega)$ we need to take care in defining the set of discrete admissible functions. A natural definition is to impose the Dirichlet condition on the face barycenters,

$$\mathcal{A}_\ell = \big\{v_\ell \in \mathrm{CR}(\mathcal{T}_\ell)^m : v_\ell^{(i)}(z) = \Pi_\ell g^{(i)}(z) \text{ for } z \in \Gamma^{(i)} \cap \mathcal{N}_\ell^{\mathrm{nc}}, i = 1, \dots, m\big\}. \tag{12}$$

We note that $\Pi_\ell v \in \mathcal{A}_\ell$, for each $v \in \mathcal{A}$, and hence, $\mathcal{A}_\ell$ is sufficiently rich to be a 'good' approximation of $\mathcal{A}$.

The Crouzeix–Raviart finite element discretization of (5) is to find a minimizer

$$u_\ell \in \operatorname{argmin} \mathcal{J}(\mathcal{A}_\ell). \tag{13}$$

As in the continuous case, we have the following result.

Proposition 3. *There exists at least one solution to* (13).

Proof. We simply adapt the proof of Proposition 1. As before, $\mathcal{A}_\ell$ is convex and closed, and $\mathcal{J}$ is (weakly) lower semicontinuous on the finite dimensional space $\mathrm{CR}(\mathcal{T}_\ell)^m$. It thus only remains to prove the coercivity of $\mathcal{J}$ in $\mathcal{A}$: a broken Poincaré inequality [8, Corollary 4.3] (for the case $p = 2$ see also [6]) reveals

$$\|v_\ell\|_{\mathrm{L}^p(\Omega)} \leq \|v_\ell - \Pi_\ell g\|_{\mathrm{L}^p(\Omega)} + \|\Pi_\ell g\|_{\mathrm{L}^p(\Omega)} \lesssim \|\nabla(v_\ell - \Pi_\ell g)\|_{\mathrm{L}^p(\Omega)} + \|\Pi_\ell g\|_{\mathrm{L}^p(\Omega)},$$

so that stability of Π_ℓ yields

$$\|v_\ell\|_{\mathrm{L}^p(\Omega)} \lesssim \|\nabla v_\ell\|_{\mathrm{L}^p(\Omega)} + \|g\|_{\mathrm{W}^{1,p}(\Omega)}.$$

With the same arguments as in the proof of Proposition 1, we thus obtain coercivity

$$\|v_\ell\|^p_{\mathrm{L}^p(\Omega)} + \|\nabla v_\ell\|^p_{\mathrm{L}^p(\Omega)} \lesssim \mathcal{J}(v_\ell) + 1 \quad \text{for all } v_\ell \in \mathcal{A}_\ell, \tag{14}$$

which is the discrete analogue of (6). Arguing as above, we conclude the proof. □

3.1. Sufficient Conditions for Convergence.

In this section, we derive conditions under which a sequence $(u_\ell)_{\ell\in\mathbb{N}}$ of discrete solutions converges to a solution of (5). Lemma 4 is the main observation which led to the convergence theorem for uniformly refined meshes [21, Equation (31)] and will again play a dominant role here. Lemma 5 is the crucial refinement of [21, Lemma 8] which will allow us to adapt the convergence argument to adaptively refined meshes.

Lemma 4 (Upper Bound). *For every $v \in \mathrm{W}^{1,p}(\Omega)^m$, it holds that*

$$\mathcal{J}(\Pi_\ell v) \leq \mathcal{J}(v) + C_{\mathrm{apx}}\,\|h_\ell f\|_{\mathrm{L}^{p'}(\Omega)}\,\|\nabla v\|_{\mathrm{L}^p(\Omega)}. \tag{15}$$

Proof. Jensen's inequality yields $W\big(|T|^{-1}\int_T \nabla v\,\mathrm{d}x\big) \leq |T|^{-1}\int_T W(\nabla v)\,\mathrm{d}x$. From the mean value property (9), we infer

$$\int_T W(\nabla \Pi_\ell v)\,\mathrm{d}x \leq \int_T W(\nabla v)\,\mathrm{d}x,$$

and hence,

$$\mathcal{J}(\Pi_\ell v) \leq \mathcal{J}(v) + \int_\Omega f\cdot(v - \Pi_\ell v)\,\mathrm{d}x.$$

Elementwise application of the approximation estimate (8) results in (15). □

Lemma 5 (Compactness of Sublevel Sets). *Let $v_\ell \in \mathcal{A}_\ell$, $\ell \in \mathbb{N}$, be a sequence satisfying $\sup_{\ell\in\mathbb{N}} \mathcal{J}(v_\ell) < \infty$, and assume that*

$$\big\|h_\ell[v_\ell]\big\|_{\mathrm{L}^1(\cup\mathcal{E}_\ell)} \xrightarrow{\ell\to\infty} 0. \tag{16}$$

(We stress that the skeleton $\bigcup \mathcal{E}_\ell$ includes $\partial\Omega$.) Then, there exists a subsequence $(v_{\ell_k})_{k\in\mathbb{N}}$ and a limit $v \in \mathcal{A}$ such that

$$\begin{aligned} v_{\ell_k} &\rightharpoonup v && \textit{weakly in } \mathrm{L}^p(\Omega)^m, \\ \nabla v_{\ell_k} &\rightharpoonup \nabla v && \textit{weakly in } \mathrm{L}^p(\Omega)^{m\times n}. \end{aligned} \tag{17}$$

Proof. As in the proof of Proposition 3 (cf. (14)) boundedness of the energy gives

$$\sup_{\ell\in\mathbb{N}} \big(\|v_\ell\|_{\mathrm{L}^p(\Omega)} + \|\nabla v_\ell\|_{\mathrm{L}^p(\Omega)}\big) < \infty.$$

Since $\mathrm{L}^p(\Omega)$ is reflexive, we may assume without loss of generality that v_ℓ as well as ∇v_ℓ are weakly convergent with limits $v \in \mathrm{L}^p(\Omega)^m$ and $F \in \mathrm{L}^p(\Omega)^{m\times n}$. We now aim to show that v is weakly differentiable with $\nabla v = F$. To this end, we fix $\varphi \in \mathrm{C}_0^\infty(\Omega)^{m\times n}$, and use (11) to obtain

$$\langle Dv_\ell, \varphi\rangle = \int_\Omega \nabla v_\ell : \varphi \,\mathrm{d}x - \int_{\cup\mathcal{E}_\ell} ([v_\ell]\otimes\nu) : \varphi\,\mathrm{d}s$$

For the first term, weak convergence of ∇v_ℓ to F implies

$$\int_\Omega \nabla v_\ell : \varphi\,\mathrm{d}x \xrightarrow{\ell\to\infty} \int_\Omega F : \varphi\,\mathrm{d}x.$$

For each face $E \in \mathcal{E}_\ell$, let $\varphi_E := |E|^{-1}\int_E \varphi\,\mathrm{d}s$ denote the integral mean of φ over E. Using the fact that $\int_E [v_\ell]\,\mathrm{d}s = 0$ for all interior faces $E \in \mathcal{E}_\ell^{\mathrm{int}}$, we estimate the second term by

$$\begin{aligned} \Big|\int_{\cup\mathcal{E}_\ell} ([v_\ell]\otimes\nu) : \varphi\,\mathrm{d}s\Big| &= \Big|\sum_{E\in\mathcal{E}_\ell}\int_E ([v_\ell]\otimes\nu) : \big(\varphi-\varphi_E\big)\,\mathrm{d}s\Big| \\ &\le \sum_{E\in\mathcal{E}_\ell} h_E \int_E \big|[v_\ell]\big|\,\mathrm{d}s\,\|\nabla\varphi\|_{\mathrm{L}^\infty(\Omega)} \\ &= \int_{\cup\mathcal{E}_\ell} h_\ell\big|[v_\ell]\big|\,\mathrm{d}s\,\|\nabla\varphi\|_{\mathrm{L}^\infty(\Omega)}. \end{aligned}$$

According to (16) and the definition of the distributional gradient, we thus obtain

$$\int_\Omega F\varphi\,\mathrm{d}x = \lim_{\ell\to\infty}\langle Dv_\ell,\varphi\rangle = -\lim_{\ell\to\infty}\int_\Omega v_\ell\ \mathrm{div}\varphi\,\mathrm{d}x = -\int_\Omega v\ \mathrm{div}\varphi\,\mathrm{d}x = \langle Dv,\varphi\rangle,$$

which proves $v \in \mathrm{W}^{1,p}(\Omega)^m$ with $\nabla v = Dv = F$.

It remains to show that $v|_{\Gamma^{(i)}} = g|_{\Gamma^{(i)}}$. Here it is crucial that (16) includes the condition that $\|h_\ell(v_\ell - g)^{(i)}\|_{\mathrm{L}^1(\Gamma^{(i)})} \to 0$. The result then follows upon combining the arguments from [21, Lemma 8] with the generalization presented above. □

Theorem 6 (Convergence of Discrete Minimizers). *Suppose that a sequence $u_\ell \in \mathrm{argmin}\mathcal{J}(\mathcal{A}_\ell)$ of discrete minimizers satisfies*

$$\|h_\ell f\|_{\mathrm{L}^{p'}(\Omega)} + \|h_\ell[u_\ell]\|_{\mathrm{L}^1(\cup\mathcal{E}_\ell)} \xrightarrow{\ell\to\infty} 0. \tag{18}$$

Then there exists a subsequence $(u_{\ell_k})_{k\in\mathbb{N}}$*, and* $u \in \operatorname{argmin} \mathcal{J}(\mathcal{A})$ *such that*

$$\begin{aligned} u_{\ell_k} &\rightharpoonup u && \text{weakly in } \mathrm{L}^p(\Omega)^m, \\ \nabla u_{\ell_k} &\rightharpoonup \nabla u && \text{weakly in } \mathrm{L}^p(\Omega)^{m\times n}, \text{ and} \\ \mathcal{J}(u_\ell) &\to \mathcal{J}(u) = \inf \mathcal{J}(\mathcal{A}). \end{aligned}$$

Moreover, unique solvability (i.e., $\#\mathrm{argmin}\mathcal{J}(\mathcal{A}) = 1$*) implies weak convergence* $u_\ell \rightharpoonup u$ *of the entire sequence. Finally, if* W *is strictly convex, it even holds that*

$$\nabla u_\ell \to \nabla u \qquad \text{strongly in } \mathrm{L}^p(\Omega)^{m\times n}.$$

Proof. Fix an arbitrary $v \in \mathcal{A}$ with finite energy. From Lemma 4, we infer

$$\mathcal{J}(u_\ell) \le \mathcal{J}(\Pi_\ell v) \le \mathcal{J}(v) + C_{\mathrm{apx}} \|h_\ell f\|_{\mathrm{L}^{p'}(\Omega)} \|\nabla v\|_{\mathrm{L}^p(\Omega)}.$$

Thus, the sequence $(u_\ell)_{\ell\in\mathbb{N}}$ has uniformly bounded energy, and hence, Lemma 5 provides a weakly convergent subsequence $(u_{\ell_k})_{k\in\mathbb{N}}$ with limit $u \in \mathcal{A}$. Since W is convex, we enjoy lower semi-continuity of $\mathcal{J}$ along the sequence u_{ℓ_k} [15, Theorem 3.4], which gives

$$\mathcal{J}(u) \le \liminf_{k\to\infty} \mathcal{J}(u_{\ell_k}) \le \limsup_{k\to\infty} \mathcal{J}(u_{\ell_k}) \le \limsup_{\ell\to\infty} \mathcal{J}(u_\ell) \le \mathcal{J}(v). \tag{19}$$

Since $v \in \mathcal{A}$ was arbitrary, we deduce $u \in \mathrm{argmin}\mathcal{J}(\mathcal{A})$. In particular, the choice $v = u$ yields equality in the latter estimate, and hence $\mathcal{J}(u) = \lim_k \mathcal{J}(u_{\ell_k}) = \inf \mathcal{J}(\mathcal{A})$ for the subsequence $(u_{\ell_k})_{k\in\mathbb{N}}$.

The convergence $\mathcal{J}(u) = \lim_\ell \mathcal{J}(u_\ell)$ follows from the fact that $\mathcal{J}(u) = \inf \mathcal{J}(\mathcal{A})$, i.e., that the limit is independent of the subsequence. Namely, if $(\widetilde{u}_\ell)_{\ell\in\mathbb{N}}$ is an arbitrary subsequence of $(u_\ell)_{\ell\in\mathbb{N}}$ for which $(\mathcal{J}(\widetilde{u}_\ell))_{\ell\in\mathbb{N}}$ is convergent, then the preceding arguments shows that $\lim_\ell \mathcal{J}(\widetilde{u}_\ell) = \inf \mathcal{J}(\mathcal{A})$. In particular, $\liminf_\ell \mathcal{J}(u_\ell) = \limsup_\ell \mathcal{J}(u_\ell) = \inf \mathcal{J}(\mathcal{A})$.

If the minimizer $u \in \mathrm{argmin}_{\mathcal{A}}\mathcal{J}$ is unique, we can use the same kind of uniqueness argument to show that the entire sequence $(u_\ell)_{\ell\in\mathbb{N}}$ converges weakly to u: More precisely, the preceding argument shows that any subsequence $(\widetilde{u}_\ell)_{\ell\in\mathbb{N}}$ of $(u_\ell)_{\ell\in\mathbb{N}}$ has a weakly convergent subsequence $(\widetilde{u}_{\ell_k})_{k\in\mathbb{N}}$, whose limit is the unique minimizer $u \in \mathcal{A}$. Consequently, the whole sequence $(u_\ell)_{\ell\in\mathbb{N}}$ converges weakly to u.

Finally, if W is strictly convex then a result of Visintin [23] shows that weak convergence together with convergence of the energy implies strong convergence. □

3.2. Refinement Indicators. The analysis of the previous section has demonstrated that condition (18) is sufficient in order to obtain convergence of the CR-FEM. It is therefore natural to use the quantities featured therein to steer the mesh refinement.

While the first convergence indicator $\|h_\ell f\|_{\mathrm{L}^{p'}(\Omega)}$ is standard and can be understood to express the approximation of the data, the second term, $\|h_\ell [u_\ell]\|_{\mathrm{L}^1(\cup\mathcal{E}_\ell)}$, is more unusual. We recall that its convergence to zero implies the weak convergence of the piecewise gradients to the gradient of a Sobolev function and hence we could interpret it as an *error estimate in the weak topology*.

In what follows, we discuss straightforward modifications of the convergence indicators which are more suitable for steering an adaptive algorithm, but for which our theory still applies. In order to associate the quantity $\|h_\ell f\|_{\mathrm{L}^{p'}(\Omega)}$ to faces $E \in \mathcal{E}_\ell$ it is natural to define a related convergence indicator as

$$\eta_\ell = \sum_{E\in\mathcal{E}_\ell} \eta_\ell(E) = \sum_{E\in\mathcal{E}_\ell} h_E^{p'} \|f\|_{\mathrm{L}^{p'}(\omega_E)}^{p'}, \tag{20}$$

where $\omega_E = \cup\{T \in \mathcal{T}_\ell : E \subset \overline{T}\}$ denotes the patch of E.

Next, applying Hölder's inequality on each face E shows

$$\begin{aligned} \sum_{E\in\mathcal{E}_\ell} \int_E h_E |[u_\ell]| \,\mathrm{d}s &\le \sum_{E\in\mathcal{E}_\ell} \Big(\int_E h_E \,\mathrm{d}s \Big)^{1/p'} \Big(\int_E h_E |[u_\ell]|^p \,\mathrm{d}s \Big)^{1/p} \\ &\lesssim |\Omega|^{1/p'} \Big(\sum_{E\in\mathcal{E}_\ell} h_E \int_E \big|[u_\ell]\big|^p \,\mathrm{d}s \Big)^{1/p}. \end{aligned} \tag{21}$$

Thus, a straightforward generalization of the indicator $\|h_\ell [u_\ell]\|_{\mathrm{L}^1(\cup\mathcal{E}_\ell)}$, is given by

$$\mu_\ell^{(0)} = \sum_{E\in\mathcal{E}_\ell} \mu_\ell^{(0)}(E) = \sum_{E\in\mathcal{E}_\ell} h_E \|[u_\ell]\|_{\mathrm{L}^p(E)}^p. \tag{22}$$

In many situations, this quantity is a bad candidate for steering the mesh refinement. To see this, let us consider the Dirichlet problem

$$-\Delta u = f \text{ in } \Omega \quad \text{with homogeneous boundary conditions} \quad u = 0 \text{ on } \partial\Omega, \tag{23}$$

where $W(F) = \frac{1}{2}|F|^2$, $p = 2$, $\Gamma^{(i)} = \partial\Omega$, and $g = 0$. For this problem, the *natural* error indicator is given by

$$\varrho_\ell^2 = \sum_{E\in\mathcal{E}_\ell\cap\Omega} h_E \int_E \big|[\nabla u_\ell]\big|^2 \,\mathrm{d}s + \sum_{E\in\mathcal{E}_\ell\cap\partial\Omega} h_E \int_E \big|\partial_\tau u_\ell\big|^2 \,\mathrm{d}s + \sum_{E\in\mathcal{E}_\ell} \eta_\ell(E),$$

where $\eta_\ell(E)$ is defined in (20) and where $\partial_\tau u_\ell$ denotes the tangential part of the gradient. This indicator is reliable and efficient (up to data oscillations) in the sense that

$$C_{\mathrm{rel}}^{-1} \|u - u_\ell\|_{\mathrm{W}^{1,2}(\Omega)} \le \varrho_\ell \le C_{\mathrm{eff}} \big(\|u - u_\ell\|_{\mathrm{W}^{1,2}(\Omega)} + \|h_\ell(f - \Pi_\ell f)\|_{\mathrm{L}^2(\Omega)} \big),$$

and it leads to a convergent adaptive algorithm [10]. Furthermore, using the fact that normal jumps can be estimated above by the $\eta_\ell(E)$ terms (cf. [10, Theorem 3.5]) and that the tangential jump of the gradient can be estimated by the tangential jump of the function, it follows easily that this indicator is equivalent to

$$\widetilde{\varrho}_\ell^2 = \sum_{E\in\mathcal{E}_\ell} \Big\{ h_E^{-1} \int_E \big|[u_\ell]\big|^2 \,\mathrm{d}s + \eta_\ell(E)^2 \Big\} = \sum_{E\in\mathcal{E}_\ell} \Big\{ h_E^{-1} \|[u_\ell]\|_{\mathrm{L}^2(E)}^2 + \eta_\ell(E)^2 \Big\}$$

This argument suggests that the indicator $\mu_\ell^{(0)}$ from (22) is not suitable, since it uses the *wrong* scaling of the mesh size. It therefore appears natural to us to use the generalization

$$\mu_\ell^{(1)} = \sum_{E\in\mathcal{E}_\ell} h_E^{1-p}\int_E \big|[u_\ell]\big|^p\,\mathrm{d}s = \sum_{E\in\mathcal{E}_\ell} h_E^{1-p}\|[u_\ell]\|^p_{\mathrm{L}^p(E)} \tag{24}$$

as an error indicator. Simple scaling arguments show why this is in fact the correct generalization (cf. [8] and Lemma 7). We will now define a further refinement indicator which can be thought of as a 'convex combination' between $\mu_\ell^{(0)}$ and $\mu_\ell^{(1)}$: for some fixed parameter $\alpha\in[0,1]$, let

$$\mu_\ell^{(\alpha)} = \sum_{E\in\mathcal{E}_\ell}\mu_\ell^{(\alpha)}(E) = \sum_{E\in\mathcal{E}_\ell} h_E^{1-\alpha p}\|[u_\ell]\|^p_{\mathrm{L}^p(E)}. \tag{25}$$

Although we have initially motivated the definition of $\mu_\ell^{(\alpha)}$ through the error estimator for the Dirichlet problem, we can give an alternative interpretation. On the Dirichlet boundary, $\mu_\ell^{(\alpha)}(E)$ weakly imposes the Dirichlet condition, while in the interior, it can be thought of as a measure of the local regularity of ∇u_ℓ. In this sense, it seems a reasonable indicator which is independent of the problem solved.

We conclude this discussion with two simple observations. The first allows us to replace the term $\|h_\ell[u_\ell]\|_{\mathrm{L}^1(\cup\mathcal{E}_\ell)}$ in Theorem 6 by $\mu_\ell^{(\alpha)}$, while the second is intended to simplify the subsequent analysis.

Lemma 7. *Suppose that $0\le\alpha\le 1$, $\ell\in\mathbb{N}$, and $u_\ell\in\mathcal{A}_\ell$, then*

$$\|h_\ell[u_\ell]\|^p_{\mathrm{L}^1(\cup\mathcal{E}_\ell)} \le C_\mu\,\mu_\ell^{(\alpha)}. \tag{26}$$

Furthermore, we have the bounds

$$\mu_\ell^{(\alpha)}(E) \le C'_\mu\|h_\ell^{1-\alpha}\nabla u_\ell\|^p_{\mathrm{L}^p(\omega_E)} \quad \text{for all interior faces } E\in\mathcal{E}_\ell^{\mathrm{int}}, \tag{27}$$

$$\mu_\ell^{(\alpha)}(E) \le C''_\mu\big(\|h_\ell^{1-\alpha}\nabla u_\ell\|^p_{\mathrm{L}^p(\omega_E)} + \|h_\ell^{1-\alpha}\nabla g\|^p_{\mathrm{L}^p(\omega_E)}\big) \quad \text{for all } E\in\mathcal{E}_\ell\backslash\mathcal{E}_\ell^{\mathrm{int}}. \tag{28}$$

All constants C_μ, C'_μ, and C''_μ depend on the shape regularity $\sigma(\mathcal{T}_\ell)$, and C_μ additionally on $|\Omega|$.

Proof. The first bound follows from (21). To prove (27), let $T^\pm\in\mathcal{T}_\ell$ denote the unique elements with $E = T^+\cap T^-\in\mathcal{E}_\ell^{\mathrm{int}}$ and $\omega_E = T^+\cup T^-$. If z_E denotes the barycentre of E, then $[u_\ell](z_E) = 0$ yields

$$|[u_\ell]| \le h_E|\nabla[u_\ell]| \le h_E(|\nabla u_\ell|_{T^+}| + |\nabla u_\ell|_{T^-}|) \quad \text{pointwise on } E.$$

Now, shape regularity of $\mathcal{T}_\ell$ gives $|T^\pm| h_E^{(1-\alpha)p} \approx |E| h_E^p h_E^{1-\alpha p}$ and results in

$$\mu_\ell^{(\alpha)}(E) \lesssim h_E^{1-\alpha p}|E| h_E^p(|\nabla u_\ell|_{T^+}|^p + |\nabla u_\ell|_{T^-}|^p) \approx h_E^{(1-\alpha)p}\|\nabla u_\ell\|^p_{\mathrm{L}^p(\omega_E)}.$$

To prove (28) note that $[u_\ell]^{(i)} = 0$ on E if $E\cap\Gamma^{(i)} = \emptyset$. We may therefore assume, without loss of generality, that $\Gamma^{(i)} = \partial\Omega$ for all $i = 1,\dots,m$, and hence $[u_\ell] = g - u_\ell$ on $\partial\Omega$. The trace inequality reads

$$h_E\|g-u_\ell\|^p_{\mathrm{L}^p(E)} \lesssim \|g-u_\ell\|^p_{\mathrm{L}^p(T)} + h_E^p\|\nabla(g-u_\ell)\|^p_{\mathrm{L}^p(T)}.$$

Note that $\int_E (g - u_\ell)\,ds = 0$ by definition of $\mathcal{A}_\ell \ni u_\ell$, and recall that in the proof of Lemma 2 it was sufficient to have mean zero on one single face to obtain the first-order approximation property. This provides $\|g - u_\ell\|_{L^p(T)} \lesssim h_T \|\nabla(g - u_\ell)\|_{L^p(T)}$, which gives

$$\mu_\ell^{(\alpha)}(E) = h_E^{1-\alpha p} \|[u_\ell]\|_{L^p(E)}^p = h_E^{1-\alpha p} \|g - u_\ell\|_{L^p(E)}^p \lesssim h_E^{p-\alpha p} \|\nabla(g - u_\ell)\|_{L^p(T)}^p$$

and immediately implies (28). □

3.3. Adaptive Strategy. We are now in a position to formulate an adaptive mesh refining strategy for the solution of (13). In what follows, $\alpha \in [0,1]$ is an arbitrary but fixed parameter of the algorithm.

Algorithm 7. **Input:** *Marking parameters $\theta \in (0,1]$, $\alpha \in [0,1]$; Initial mesh $\mathcal{T}_0$. Set $\ell = 0$.*

(a) *Compute a discrete minimizer $u_\ell \in \operatorname{argmin} \mathcal{J}(\mathcal{A}_\ell)$.*
(b) *Compute refinement indicators η_ℓ and $\mu_\ell^{(\alpha)}$ from* (20) *and* (25)*, respectively.*
(c) *Generate a set of marked faces $\mathcal{M}_\ell \subseteq \mathcal{E}_\ell$ such that*

$$\sum_{E \in \mathcal{M}_\ell} \big(\eta_\ell(E) + \mu_\ell^{(\alpha)}(E)\big) \geq \theta(\eta_\ell + \mu_\ell^{(\alpha)}) \tag{29}$$

(d) *Generate a regular triangulation $\mathcal{T}_{\ell+1}$, where at least the marked faces $E \in \mathcal{M}_\ell$ are refined.*
(e) *Increase $\ell \mapsto \ell + 1$ and go to* (a). □

A marking strategy satisfying (29) is often called *Dörfler marking.* It was a crucial ingredient in the first convergence proofs of the adaptive finite element method and has been identified to also play an important role in obtaining optimal convergence rates [13], in which case the cardinality of the set $\mathcal{M}_\ell$ should be minimal. Generically, the value $\theta = 1$ corresponds to uniform refinement, whereas small θ leads to highly adapted meshes.

We use newest vertex bisection in step (4) to ensure that the sequence of triangulations $(\mathcal{T}_\ell)_{\ell \in \mathbb{N}}$ generated by Algorithm 1 is uniformly shape-regular. However, besides the uniform shape-regularity, the following convergence result only requires that marked faces are reduced by a uniform factor $\kappa \in (0,1)$, i.e.,

$$h_{E'} \leq \kappa\, h_E \quad \text{for all } E' \in \mathcal{E}_{\ell+1} \text{ with } E' \subset E \in \mathcal{M}_\ell. \tag{30}$$

Thus, the precise refinement rule in step (d) is fairly arbitrary.

Theorem 8 (Convergence of Adaptive Algorithm). *Suppose that the sequence $(\mathcal{T}_\ell)_{\ell \in \mathbb{N}}$ generated by Algorithm 1 is uniformly shape-regular, that it satisfies* (30)*, and assume in addition that $0 \leq \alpha < 1$. Then the refinement indicators converge to zero, i.e.,*

$$\eta_\ell + \mu_\ell^{(\alpha)} \xrightarrow{\ell \to \infty} 0, \tag{31}$$

and the sequence $(u_\ell)_{\ell \in \mathbb{N}}$ of discrete minimizers satisfies the conditions of Theorem 6.

Proof. As in the proof of Theorem 6, it follows that $\mathcal{J}(u_\ell)$ and hence $\|u_\ell\|_{\mathrm{L}^p(\Omega)} + \|\nabla u_\ell\|_{\mathrm{L}^p(\Omega)}$ are bounded sequences.

To abbreviate the notation, we now drop the superscript (α) in $\mu_\ell^{(\alpha)}$, and we write

$$\eta_\ell(\mathcal{S}_\ell) := \sum_{E\in\mathcal{S}_\ell} \eta_\ell(E) \quad \text{and} \quad \mu_\ell(\mathcal{S}_\ell) := \sum_{E\in\mathcal{S}_\ell} \mu_\ell(E) \quad \text{for all } \mathcal{S}_\ell \subseteq \mathcal{E}_\ell.$$

We consider the set of all faces resp. all elements which are eventually not refined, i.e.,

$$\widetilde{\mathcal{E}} := \bigcap_{k\geq 0}\bigcup_{\ell\geq k} \mathcal{E}_\ell \quad \text{and} \quad \widetilde{\mathcal{T}} := \bigcap_{k\geq 0}\bigcup_{\ell\geq k} \mathcal{T}_\ell.$$

It is evident that $T \in \widetilde{\mathcal{T}}$ if and only if all faces of T belong to $\widetilde{\mathcal{E}}$. For the proof of (31), we split the indicators into

$$\eta_\ell + \mu_\ell = \big(\eta_\ell(\mathcal{E}_\ell\backslash\widetilde{\mathcal{E}}) + \mu_\ell(\mathcal{E}_\ell\backslash\widetilde{\mathcal{E}})\big) + \big(\eta_\ell(\mathcal{E}_\ell\cap\widetilde{\mathcal{E}}) + \mu_\ell(\mathcal{E}_\ell\cap\widetilde{\mathcal{E}})\big)$$

Step 1: In the first step, we will prove that

$$\eta_\ell(\mathcal{E}_\ell \setminus \widetilde{\mathcal{E}}) + \mu_\ell(\mathcal{E}_\ell \setminus \widetilde{\mathcal{E}}) \xrightarrow{\ell\to\infty} 0. \tag{32}$$

Recall that $\omega_E := \bigcup\big\{T \in \mathcal{T}_\ell \,:\, E \subset \partial T\big\} \subseteq \overline{\Omega}$ denotes the patch of E. With $\widetilde{\Omega}_\ell = \bigcup\big\{\omega_E \,:\, E \in \mathcal{E}_\ell \setminus \widetilde{\mathcal{E}}\big\}$, we first claim that $\chi_{\widetilde{\Omega}_\ell} h_\ell \xrightarrow{\ell\to\infty} 0$ almost everywhere in Ω. To see this, fix $x \in \Omega\backslash\big(\bigcup_\ell \bigcup \mathcal{E}_\ell\big)$ outside of the skeletons of all $\mathcal{T}_\ell$ which form a null-set. For each ℓ, there is a unique element $T_\ell \in \mathcal{T}_\ell$ with $x \in T_\ell$. If $\lim_\ell h_{T_\ell} = 0$, we conclude $\lim_\ell(\chi_{\widetilde{\Omega}_\ell} h_\ell)(x) = 0$. Otherwise, T_ℓ is only refined finitely many times, i.e., there holds $T_\ell = T_{\ell_0}$ for some $\ell_0 \in \mathbb{N}$ and all $\ell \geq \ell_0$, i.e., $T_\ell \in \widetilde{\mathcal{T}}$ and therefore its faces belong to $\widetilde{\mathcal{E}}$. Consequently, $x \notin \widetilde{\Omega}_\ell$ for all $\ell \geq \ell_0$, and hence $(\chi_{\widetilde{\Omega}_\ell} h_\ell)(x) = 0$ for $\ell \geq \ell_0$. We have therefore shown that

$$\lim_{\ell\to\infty} \chi_{\widetilde{\Omega}_\ell} h_\ell = 0 \qquad \text{pointwise a.e. in } \Omega.$$

During mesh-refinement, the local mesh-size h_ℓ is pointwise decreasing. Consequently, the dominated convergence theorem yields

$$\chi_{\widetilde{\Omega}_\ell} h_\ell^\beta \psi \xrightarrow{\ell\to\infty} 0 \qquad \text{strongly in } \mathrm{L}^q(\Omega), \tag{33}$$

for all $\beta > 0$, and $\psi \in \mathrm{L}^1(\Omega)$. With $\beta = 1$ and $\psi = f$, we infer

$$\eta_\ell(\mathcal{E}_\ell\backslash\widetilde{\mathcal{E}}) = \sum_{E\in\mathcal{E}_\ell\backslash\widetilde{\mathcal{E}}} \|h_\ell f\|_{\mathrm{L}^{p'}(\omega_E)}^{p'} \lesssim \|h_\ell f\|_{\mathrm{L}^{p'}(\widetilde{\Omega}_\ell)}^{p'} = \|\chi_{\widetilde{\Omega}_\ell} h_\ell f\|_{\mathrm{L}^{p'}(\Omega)}^{p'} \xrightarrow{\ell\to\infty} 0.$$

Before we prove convergence of $\mu_\ell(\mathcal{E}_\ell\backslash\widetilde{\mathcal{E}})$ to zero, it is instructive to consider the refinement indicator $\|h_\ell[u_\ell]\|_{\mathrm{L}^1(\cup\mathcal{E}_\ell)}$ first. Using the facts that $h_E \approx h_\ell$ in ω_E, and that $[u_\ell](z_E) = 0$, we can estimate

$$\int_E h_E|[u_\ell]|\,\mathrm{d}s \lesssim h_E^2 \int_E \big|[\nabla u_\ell]\big|\,\mathrm{d}s \lesssim \int_{\omega_E} h_\ell|\nabla u_\ell|\,\mathrm{d}x.$$

Summing over $E \in \mathcal{E}_\ell \setminus \widetilde{\mathcal{E}}$ and using Hölder's inequality, we obtain

$$\|h_\ell[u_\ell]\|_{\mathrm{L}^1(\cup(\mathcal{E}_\ell\setminus\widetilde{\mathcal{E}}))} \lesssim \sum_{E\in\mathcal{E}_\ell\setminus\widetilde{\mathcal{E}}} \|h_\ell \nabla u_\ell\|_{\mathrm{L}^1(\omega_E)} \lesssim \|h_\ell \nabla u_\ell\|_{\mathrm{L}^1(\widetilde{\Omega}_\ell)} \le \|h_\ell\|_{\mathrm{L}^{p'}(\widetilde{\Omega}_\ell)} \|\nabla u_\ell\|_{\mathrm{L}^p(\widetilde{\Omega}_\ell)}.$$

According to (33), with $\beta = 1$ and $\psi = 1$, the upper bound tends to zero as $\ell \to \infty$.

If we attempt to use the same idea for proving that $\mu_\ell(\mathcal{E}_\ell \setminus \widetilde{\mathcal{E}}) \to 0$ then, using (27) and (28), we first obtain the following bound:

$$\mu_\ell(\mathcal{E}_\ell \setminus \widetilde{\mathcal{E}}) \lesssim \|h_\ell^{(1-\alpha)} \nabla u_\ell\|^p_{\mathrm{L}^p(\widetilde{\Omega}_\ell)} + \|h_\ell^{(1-\alpha)} \nabla g\|^p_{\mathrm{L}^p(\widetilde{\Omega}_\ell)}. \tag{34}$$

Using (33), with $\beta = 1 - \alpha > 0$ and $\psi = \nabla g$, we immediately find that the second integral on the right-hand side of (34) converges to zero. It thus only remains to to deal with the first term on the right-hand side. Unfortunately, we control ∇u_ℓ only in $\mathrm{L}^p(\Omega)^{m\times n}$. Therefore, we cannot immediately use Hölder's inequality as before to verify that the first term on the right-hand side of (34) tends to zero. Since we don't know whether ∇u_ℓ converges pointwise a.e., we have no hope of using Fatou's lemma either. Instead, we make use of the additional flexibility provided by the condition $\alpha < 1$ to estimate

$$\begin{aligned}
\|h_\ell^{(1-\alpha)} \nabla u_\ell\|^p_{\mathrm{L}^p(\widetilde{\Omega}_\ell)} = \int_{\widetilde{\Omega}_\ell} h_\ell^{(1-\alpha)p} |\nabla u_\ell|^p \,\mathrm{d}x &\lesssim \sum_{E\in\mathcal{E}_\ell\setminus\widetilde{\mathcal{E}}} \int_{\omega_E} h_\ell^{(1-\alpha)p} |\nabla u_\ell|^p \,\mathrm{d}x \\
&\lesssim \sum_{T\in\mathcal{T}_\ell\cap\widetilde{\Omega}_\ell} h_T^{n+(1-\alpha)p} |\nabla u_\ell|_T|^p \\
&\le \Big(\sum_{T\in\mathcal{T}_\ell\cap\widetilde{\Omega}_\ell} h_T^{(n+(1-\alpha)p)q/p} |\nabla u_\ell|_T|^q \Big)^{p/q} \\
&\lesssim \Big(\sum_{T\in\mathcal{T}_\ell\cap\widetilde{\Omega}_\ell} h_T^{nq/p+(1-\alpha)q-n} \int_T |\nabla u_\ell|^q \Big)^{p/q},
\end{aligned}$$

where $1 \le q < p$. In the third estimate above we used the bound $\|\cdot\|_{\ell^p} \le \|\cdot\|_{\ell^q}$. Setting $\beta = n(q/p - 1) + (1-\alpha)q$, we obtain

$$\|h_\ell^{(1-\alpha)} \nabla u_\ell\|^p_{\mathrm{L}^p(\widetilde{\Omega})} \lesssim \Big(\int_{\widetilde{\Omega}_\ell} h_\ell^\beta |\nabla u_\ell|^q \,\mathrm{d}x \Big)^{p/q} \le \|h_\ell^\beta \chi_{\widetilde{\Omega}_\ell}\|^{p/q}_{\mathrm{L}^{p/(p-q)}} \|\nabla u_\ell\|^p_{\mathrm{L}^p(\Omega)}.$$

by use of Hölder's inequality. For this bound to tend to zero, we require that $\beta > 0$ which can be achieved by choosing q sufficiently close to p and using the fact that $\alpha < 1$. Thus, we have successfully established (32).

Step 2: In the second step, we use the properties of our marking strategy to conclude the proof of (31). Namely, observe that at step ℓ any marked face will be refined during this step, i.e.,

$$\mathcal{M}_\ell \subset \mathcal{E}_\ell \setminus \widetilde{\mathcal{E}}.$$

Therefore, (29) and *Step 1* imply that

$$\theta(\eta_\ell + \mu_\ell) \le \eta_\ell(\mathcal{M}_\ell) + \mu_\ell(\mathcal{M}_\ell) \le \eta_\ell(\mathcal{E}_\ell \setminus \widetilde{\mathcal{E}}) + \mu_\ell(\mathcal{E}_\ell \setminus \widetilde{\mathcal{E}}) \to 0,$$

as $\ell \to \infty$, which concludes the proof. □

Remark 1. We have already remarked in Section 3.2 that, for $\alpha < 1$, our refinement indicators are not reliable error indicators, even for a simple Dirichlet problem with homogeneous boundary conditions. Furthermore, our refinement indicators have no information about the free boundary which gives further indication to its "incompleteness". We found it therefore somewhat surprising that we were able to prove convergence of our adaptive strategy.

We note, however, that our proof does not extend in an obvious way to conforming finite element methods, where the strong upper bound (15) is unavailable even for quadratic functionals. □

Remark 2. We have only proven convergence of our adaptive algorithm for the case $\alpha < 1$. It does not appear straightforward to include the case $\alpha = 1$ as well. The analysis in [20] shows that obtaining strong convergence of the sequence $(u_\ell)_{\ell\in\mathbb{N}}$ to some $\widetilde{u}$ *a priori* (instead of merely weak convergence) is the key. However, this appears difficult for problems of the generality which we consider here.

In practice, one may safely ignore this fact and choose $\alpha = 1$, possibly implementing a safeguard strategy which changes α if it should become apparent that $\eta_\ell + \mu_\ell \not\to 0$. □

4. Numerical Experiments

We have implemented Algorithm 1 for two two-dimensional model problems: the Laplace problem with Dirichlet and with Neumann boundary conditions as well as the example of Foss, Hrusa, and Mizel [18] which exhibits a Lavrentiev gap. Before we present the computational experiments, we briefly outline the details of our implementation.

(a) The solution of the optimization problem is achieved by a damped Newton method if it is nonlinear and a direct solver if it is linear.

(b) We have found that Dörfler's marking strategy with a minimal set $\mathcal{M}_\ell$ yields very slow mesh growth for the highly nonlinear and singular problems which we consider here. Therefore, our strategy is to mark a fixed fraction of edges (with largest indicators) for refinement, i.e.,

$$\sharp\mathcal{M}_\ell \geq \theta\sharp\mathcal{E}_\ell \quad \text{with} \quad \min_{E\in\mathcal{M}_\ell}\big(\eta_\ell(E) + \mu_\ell^{(\alpha)}(E)\big) \geq \max_{E\in\mathcal{E}_\ell\setminus\mathcal{M}_\ell}\big(\eta_\ell(E) + \mu_\ell^{(\alpha)}(E)\big).$$

Note that then,

$$\sum_{E\in\mathcal{M}_\ell}\big(\eta_\ell(E) + \mu_\ell^{(\alpha)}(E)\big) \geq \frac{\theta}{1+\theta}\sum_{E\in\mathcal{E}_\ell}\big(\eta_\ell(E) + \mu_\ell^{(\alpha)}(E)\big),$$

so that Dörfler marking (29) still holds with θ replaced by $\theta/(1+\theta)$. We usually chose $\theta = 0.25$ which roughly doubles the number of elements.

(c) The mesh refinement is achieved via *newest vertex bisection* which halves every marked edge and which preserves shape regularity.

(d) In practice, we terminate the algorithm when a prescribed number of elements is attained.

For the computations in Section 4.1 we estimate the error for the energy by comparing it to a conforming computation. If $f \equiv 0$, then

$$\mathcal{J}(u_\ell) \le \inf \mathcal{J}(\mathcal{A}) \le \mathcal{J}(\bar{u}) \quad \text{for all } \bar{u} \in \mathcal{A}.$$

If f is non-zero then the above estimate depends on an unknown quantity, namely $\|\nabla u\|_{\mathrm{L}^p}$ where $u \in \mathrm{argmin}\mathcal{J}(\mathcal{A})$. However, we have observed that even in that case, $\mathcal{J}(u_\ell)$ is monotonically increasing towards the energy of the exact solution. Therefore, we compute a conforming $\bar{u}$ using a standard adaptive P_1-finite element method [9] and take

$$\inf \mathcal{J}(\mathcal{A}) - \mathcal{J}(u_\ell) \le \mathcal{J}(\bar{u}) - \mathcal{J}(u_\ell)$$

as a slightly heuristic error estimate.

4.1. Linear Laplacian. We begin our experiments with the Laplace equation on the slit domain $\Omega_s = (-1,1)^2 \setminus [0,1) \times \{0\}$. It is equivalently formulated by setting $W(F) = \frac{1}{2}|F|^2$ with $m = 1$ and $n = 2$. First, we consider the pure Dirichlet problem

$$-\Delta u = 1 \text{ in } \Omega \quad \text{with homogeneous boundary conditions} \quad u = 0 \text{ on } \partial\Omega,$$

where in the energy formulation

$$\Gamma^{(1)} = \partial\Omega, \quad f = 1, \quad g = 0. \tag{35}$$

In order to investigate the effect of a Neumann boundary, we also consider the mixed boundary value problem

$$-\Delta u = 0 \text{ in } \Omega \quad \text{with} \quad \partial u/\partial n = 0 \text{ on } \Gamma_N \quad \text{and} \quad u = g \text{ on } \Gamma_D,$$

where $|\Gamma_D \cap \Gamma_N| = 0$ and $\partial\Omega = \Gamma_D \cup \Gamma_N$. We choose

$$\Gamma^{(1)} = \Gamma_D = \partial\Omega \cap \{x_1 = 1\}, \quad f = 0, \quad g(1, x_2) = \mathrm{sign}(x_2). \tag{36}$$

The convergence rates for problems (35) and (36) are shown in Figures 2 and 3, respectively. As expected, we observe that the accuracy improves as α approaches 1.0. In the Dirichlet problem, the convergence rate for $\alpha = 1.0$ and for $\alpha = 0.9$ can barely be distinguished. What is surprising though is that, for the Neumann problem, the value of α does not seem to affect the convergence rate at all. We have no explanation for this effect, however, we will observe it again in later examples.

4.2. Lavrentiev Phenomenon. We conclude our numerical experiments with an example which exhibits the Lavrentiev gap phenomenon — the focus of our investigation. To this end, we slightly modify the example of Foss, Hrusa and Mizel [18]. Let $m = n = 2$, let Ω be the half disk $\Omega = \{|x| < 1, x_2 > 0\}$, and let

$$\Gamma^{(1)} = (-1,0) \times \{0\} \cup \{|x| = 1, x_2 > 0\} \quad \text{and} \quad \Gamma^{(2)} = (0,1) \times \{0\} \cup \{|x| = 1, x_2 > 0\}.$$

Furthermore, $f \equiv 0$, $g^{(i)} = 0$ on $\{x_2 = 0\}$, and $g = (\cos\frac{\theta}{2}, \sin\frac{\theta}{2})$ on $\{|x| = 1\}$ in polar coordinates (r, θ). Thus, admissible functions are deformations of the half disk

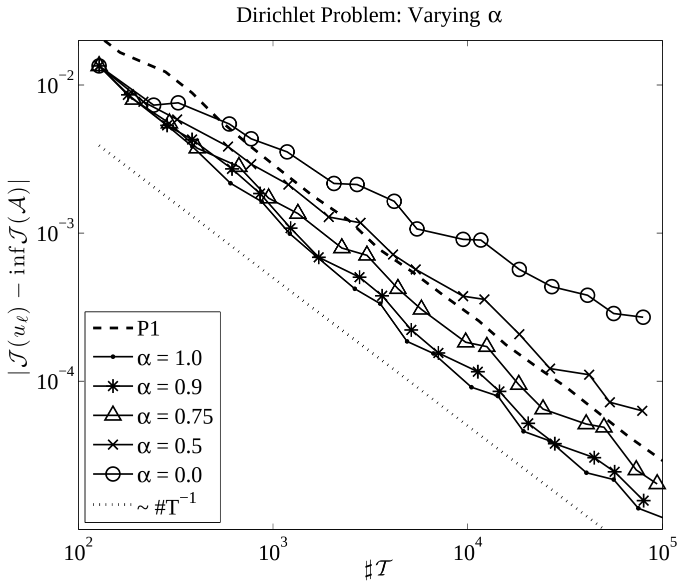

FIGURE 2. Convergence rates for Algorithm 1 applied to problem (35). As $\alpha \nearrow 1$, the convergence rate approaches the optimal rate $\sharp\mathcal{T}^{-1}$.

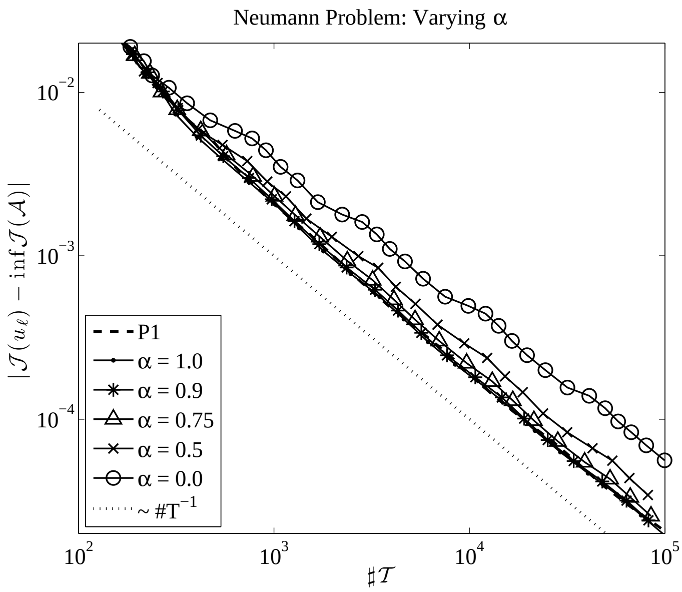

FIGURE 3. Convergence rates for Algorithm 1 applied to problem (36). Contrary to intuition, the convergence rate seems to be optimal, independent of the parameter α.

Ω into the quarter disk $\{|x| < 1,\ x_1 > 0,\ x_2 > 0\}$. Suppose, for the moment, that the stored energy density is given by

$$W(F) = \big(|F|^2 - 2\det F\big)^4.$$

Convexity of W follows immediately from the fact that $F \mapsto (|F|^2 - 2\det F)$ is a non-negative quadratic form.

The associated energy functional is not coercive. Nevertheless, Foss, Hrusa, and Mizel showed in [18] that it exhibits the Lavrentiev gap phenomenon. The global minimum of $\mathcal{J}$ in $\mathcal{A}$, with $W = W_0$, is the function

$$\bar{u} = r^{1/2}\big(\cos\tfrac{\theta}{2}, \sin\tfrac{\theta}{2}\big),$$

for which $\mathcal{J}(\bar{u}) = 0$. It is furthermore easy to verify that $\bar{u}$ also minimizes the Dirichlet integral for the same boundary conditions. Consequentlym, for $p = 2$, $\bar{u}$ is also the global minimizer of

$$\mathcal{J}_p(v) = \int_\Omega W_p(\nabla v)\,\mathrm{d}x, \tag{37}$$

in $\mathcal{A}$, where

$$W_p(F) = \big(|F|^2 - 2\det F\big)^4 + \tfrac{1}{p}\big(|F_1|^p + |F_2|^p\big).$$

Since we know the solution for $p = 2$ explicitly, we know the energy

$$\inf \mathcal{J}_2(\mathcal{A}) = \mathcal{J}_2(\bar{u}) = \pi/4.$$

In Figure 4, we plot the convergence rate for the minimization problem, for varying α. We observe the same effect as for the Neumann problem in Section 4.1: surprisingly, the convergence rate seems to be independent of the parameter α. It is also particularly encouraging that the convergence rate for the energy appears to be linear, despite the fact that we are solving a highly non-linear and singular problem. Note that $\mathcal{J}_2$ is not even continuous in the strong topology of $\mathrm{W}^{1,2}(\Omega)^2$.

4.3. Verification of Lavrentiev gaps. In our final experiment, we demonstrate how one could verify whether a given minimization problem exhibits a Lavrentiev gap. We consider the energy function $\mathcal{J}_p$ from (37). For certain fixed parameters $p = 2, 3, 4, 6$, we apply Algorithm 1 with the minimization problem $\mathrm{argmin}\mathcal{J}_p(\mathcal{A})$ and obtain discrete solutions u_ℓ. In addition, we compute an adaptive P_1-solution $\bar{u}_\ell$ for the same problem (though possibly on different meshes) and we plot the difference in energy $\mathcal{J}_p(\bar{u}_\ell) - \mathcal{J}_p(u_\ell)$. The theory in [18] would lead us to expect (but except for the case $p = 2$ this is not at all clear) that, for $p = 2, 3$ a Lavrentiev gap occurs, while for $p = 4, 6$ no gap is present. The computations which we show in Figure 5 fully agree with this prediction, except in the case $p = 4$, where they suggest that no Lavrentiev gap may in fact be present.

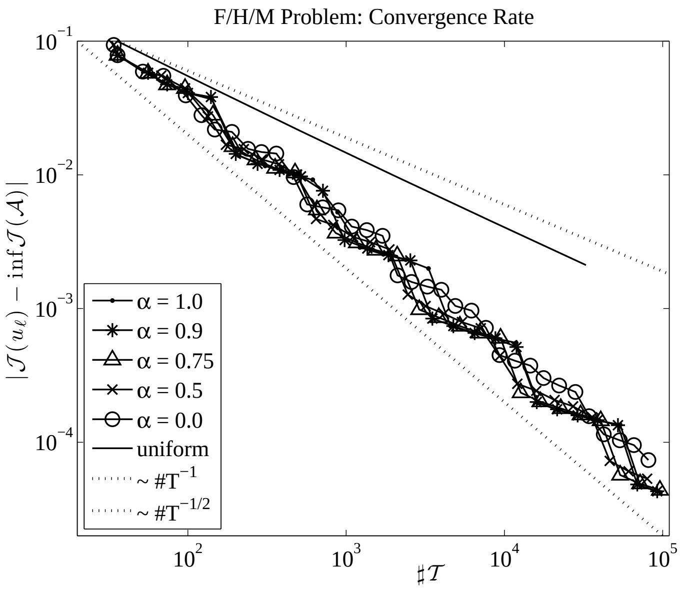

Figure 4. Convergence rates for Algorithm 1 applied to the minimization problem $u \in \operatorname{argmin} \mathcal{J}_2(\mathcal{A})$, with varying marking parameter α.

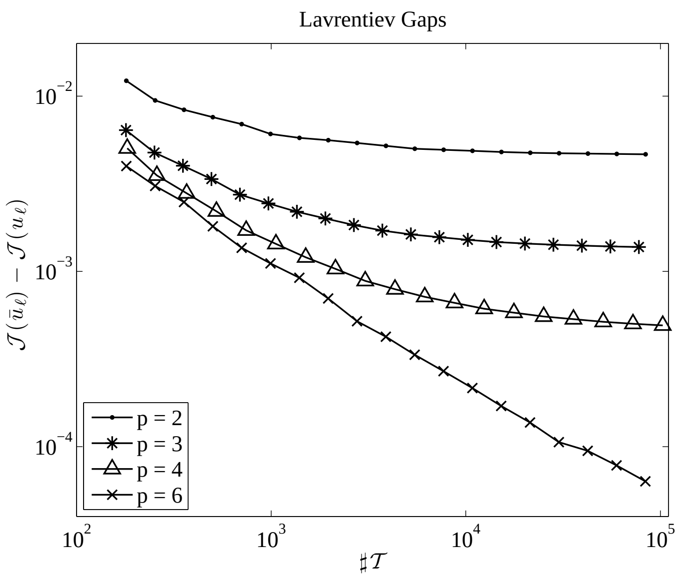

Figure 5. Adaptive computation of the Lavrentiev gap $\inf \mathcal{J}_p(\mathcal{A}_\infty) - \inf \mathcal{J}_p(\mathcal{A})$ for $p = 2, 3, 4, 6$. Contrary to intuition, for $p = 4$, the computation suggests that a Lavrentiev gap is present.

5. Conclusion

We have presented an adaptive finite element algorithm for the solution of convex variational problems. Despite the fact that the refinement indicators are *not* reliable error indicators in any classical sense, we have succeeded in proving convergence of our adaptive scheme. To conclude, we briefly mention some possible generalizations of our analysis.

In order for the minimization problem (5) to be well-posed (or rather, for the direct method technique to apply), it is necessary that $\mathcal{J}$ is coercive in $\mathcal{A}$ which, in particular, requires that elements of $\mathcal{A}$ satisfy a Poincaré-type inequality. Similarly, we require a broken Poincaré-type inequality for the discrete admissible set $\mathcal{A}_\ell$ in order to be able to extract weakly convergent subsequences. The entire analysis applies whenever such a broken Poincaré-type inequality is available for $\mathcal{A}_\ell$. It is therefore straightforward to generalize the results to obstacle problems, plasticity, and so forth.

A second important generalization is to allow W to depend on x and on u. It is not at all clear in which generality this can be achieved. Mild dependencies such as piecewise constant dependence on x are easily included in the analysis, however, a strong coupling of (x, u) to ∇u must be avoided. This follows immediately upon considering Manià's functional

$$\mathcal{J}(u) = \int_0^1 u_x^6 (u^3 - x)^2 \, \mathrm{d}x,$$

which is to be minimized subject to $u(0) = 0, u(1) = 1$. Since the CR-FEM reduces to the P_1-FEM in one dimension, and since Manià's example exhibits a Lavrentiev gap, it follows that a sufficiently strong coupling of (x, u) to ∇u destroys the convergence of the method.

Finally, the generalization to polyconvex or even quasiconvex W is even more difficult. Here, both the upper bound (15) and the lower bound (19), i.e., the weak lower-semicontinuity of $\mathcal{J}$ along the sequence u_ℓ, are completely open.

References

[1] G. Acosta and R. G. Durán. An optimal Poincaré inequality in L^1 for convex domains. *Proc. Amer. Math. Soc.*, 132(1):195–202 (electronic), 2004.

[2] R.A. Adams and J.F. Fournier. *Sobolev Spaces*. Elsevier Science, 2003.

[3] J. M. Ball. Singularities and computation of minimizers for variational problems. In *Foundations of computational mathematics (Oxford, 1999)*, volume 284 of *London Math. Soc. Lecture Note Ser.*, pages 1–20. Cambridge Univ. Press, Cambridge, 2001.

[4] S. Bartels and C. Carstensen. A convergent adaptive finite element method for an optimal design problem. *Numer. Math.*, 108(3):359–385, 2008.

[5] D. Braess. *Finite Elements*. Cambridge University Press, Cambridge, third edition, 2007. Theory, fast solvers, and applications in elasticity theory, Translated from the German by Larry L. Schumaker.

[6] Susanne C. Brenner. Poincaré-Friedrichs inequalities for piecewise H^1 functions. *SIAM J. Numer. Anal.*, 41(1):306–324 (electronic), 2003.

[7] F. Brezzi and M. Fortin. *Mixed and hybrid finite element methods*, volume 15 of *Springer Series in Computational Mathematics*. Springer-Verlag, New York, 1991.

[8] A. Buffa and C. Ortner. Compact embeddings of broken Sobolev Spaces and applications. to appear in IMA J. Numer. Anal.

[9] C. Carstensen. Convergence of an adaptive FEM for a class of degenerate convex minimization problems. to appear in IMA J. Numer. Anal.

[10] C. Carstensen and R. H. W. Hoppe. Convergence analysis of an adaptive nonconforming finite element method. *Numer. Math.*, 103(2):251–266, 2006.

[11] C. Carstensen and R. H. W. Hoppe. Error reduction and convergence for an adaptive mixed finite element method. *Math. Comp.*, 75(255):1033–1042 (electronic), 2006.

[12] C. Carstensen and C. Ortner. Gamma convergence in computational calculus of variations. in preparation.

[13] J.M. Cascon, C. Kreuzer, R.H. Nochetto, and K. G. Siebert. Quasi-optimal convergence rate for an adaptive finite element method. Technical Report 009/2007, Institut für Mathematik, Universität Augsburg, 2007.

[14] M. Crouzeix and P.-A. Raviart. Conforming and nonconforming finite element methods for solving the stationary Stokes equations. I. *Rev. Française Automat. Informat. Recherche Opérationnelle Sér. Rouge*, 7(R-3):33–75, 1973.

[15] B. Dacorogna. *Direct methods in the calculus of variations*, volume 78 of *Applied Mathematical Sciences*. Springer-Verlag, Berlin, 1989.

[16] L. Diening and C. Kreuzer. Linear convergence of an adaptive finite element method for the p-Laplacian equation. *SIAM J. Numer. Anal.*, 46(2):614–638, 2008.

[17] W. Dörfler. A convergent adaptive algorithm for Poisson's equation. *SIAM J. Numer. Anal.*, 33(3):1106–1124, 1996.

[18] M. Foss, W. J. Hrusa, and V. J. Mizel. The Lavrentiev gap phenomenon in nonlinear elasticity. *Arch. Ration. Mech. Anal.*, 167(4):337–365, 2003.

[19] P. Morin, R. H. Nochetto, and K. G. Siebert. Data oscillation and convergence of adaptive FEM. *SIAM J. Numer. Anal.*, 38(2):466–488 (electronic), 2000.

[20] P. Morin, K. Siebert, and A. Veeser. A basic convergence result for conforming adaptive finite elements. Technical Report 007/2007, Institut für Mathematik, Universität Augsburg, 2007.

[21] C. Ortner. A non-conforming finite element method for convex minimization problems. submitted for publication.

[22] A. Veeser. Convergent adaptive finite elements for the nonlinear Laplacian. *Numer. Math.*, 92(4):743–770, 2002.

[23] A. Visintin. Strong convergence results related to strict convexity. *Comm. Partial Differential Equations*, 9(5):439–466, 1984.